# How to calculate a three-dimensional $g(r)$ under periodic boundary conditions 

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(Dated: March 2, 2004)


#### Abstract

If under simple periodic boundary conditions one wants to compute the pair distribution function $g(r)$ beyond half the simple cubic box length, minimum image issues require an extension of the standard procedure. These notes summarize the required modifications.


The pair correlation function $g(r)$ may be defined as the probability density of finding a particle a distance $r$ away from another particle, divided by the probability density for the same event in a noninteracting system. Equivalently, we can take the number $N(r)$ of particles within a distance range between $r$ and $r+\mathrm{d} r$ and divide it by the number $N_{\mathrm{id}}(r)$ within the same interval in an ideal system. If $\rho$ is the average particle density, we evidently have

$$
\begin{equation*}
N_{\mathrm{id}}(r)=\rho \times 4 \pi r^{2} \mathrm{~d} r \tag{1}
\end{equation*}
$$

In other words, apart from a trivial prefactor $\rho$, we essentially divide $N(r)$ by the surface area of a sphere of radius $r$. This of course takes care of the fact that the farther away we go from the central particle, the more possibilities one has to put another particle at a distance $r$.

However, this simple reasoning does not work under periodic boundary conditions. Or, more precisely, it only works provided that $r \leq \frac{1}{2} L$ - where we assume that we have a cubic box with box length $L$. The reason is quite obvious: Once $r$ becomes bigger than $\frac{1}{2} L$, there are certain directions along which the minimum image distance between two particles is shorter than $r$. Or, in other words, even for an ideal system the number of particles at a distance $r$ does not keep increasing like $4 \pi r^{2}$.

There are two solutions to this problem: A plain one and a clever one. The plain one is: don't plot $g(r)$ beyond $r=\frac{1}{2} L$. This is acceptable, but may be a waste of information.

The clever solution is to find out, what the actual number of particles at a given minimum image distance $r$ is. This basically can be reduced to the following problem:

Given a cubic simulation box of side length 1, what is the probability density $p(r)$ for having a particular minimum image distance $r$ ?

The answer to this question is

$$
p(r)=\left\{\begin{array}{ccc}
4 \pi r^{2} & , & 0<2 r \leq 1  \tag{2}\\
2 \pi r(3-4 r) & , & 1<2 r \leq \sqrt{2} \\
2 r\left[3 \pi-12 f_{1}(r)+f_{2}(r)\right] & , \sqrt{2}<2 r \leq \sqrt{3}
\end{array},\right.
$$

where we used the abbreviations

$$
\begin{align*}
& f_{1}(r)=\arctan \sqrt{4 r^{2}-2}  \tag{3a}\\
& f_{2}(r)=8 r \arctan \frac{2 r\left(4 r^{2}-3\right)}{\sqrt{4 r^{2}-2}\left(4 r^{2}+1\right)} \tag{3b}
\end{align*}
$$



FIG. 1: Probability density $p(r)$ for the minimum image distance between two points randomly picked from the unit cube-Eqn. (2).

The graph of $p(r)$ is sketched in Fig. 1. The required modification now is simply to divide the number of particles $N(r)$ within a certain minimum image distance between $r$ and $r+\mathrm{d} r$ by the following modified normalization number:

$$
\begin{equation*}
\tilde{N}_{\mathrm{id}}(r)=\rho \times p(r / L) L^{2} \mathrm{~d} r \tag{4}
\end{equation*}
$$

Sometimes it is useful to also know the integrated probability density, i.e., the probability distribution

$$
\begin{equation*}
D(r):=\int_{0}^{r} \mathrm{~d} \bar{r} p(\bar{r}) \tag{5}
\end{equation*}
$$

for instance if bins are wide and one needs a better expression for their volume $V=D(r+\Delta r)-D(r)=p(r) \Delta r+\mathcal{O}\left(\Delta r^{2}\right)$. Luckily, the integration can be done analytically:

$$
D(r)=\left\{\begin{array}{cc}
\frac{4 \pi}{3} r^{3} & , \quad 0<2 r \leq 1  \tag{6}\\
-\frac{\pi}{12}\left(3-36 r^{2}+32 r^{3}\right) & , \quad 1<2 r \leq \sqrt{2} \\
-\frac{\pi}{4}+3 \pi r^{2}+\sqrt{4 r^{2}-2}+ & \sqrt{2}<2 r \leq \sqrt{3} \\
\left(1-12 r^{2}\right) f_{1}(r)+\frac{2}{3} r^{2} f_{2}(r) & , \quad
\end{array}\right.
$$

After these modifications $g(r)$ can be plotted up to $\frac{1}{2} \sqrt{3} L$, respectable $73 \%$ farther than the trivial range. Beyond this value it finally ceases to exist, because there simply is no larger minimum image separation available in a cubic box.

